



***Energy dissipation of relativistic outflow***  
***~Internal and External Shocks~***

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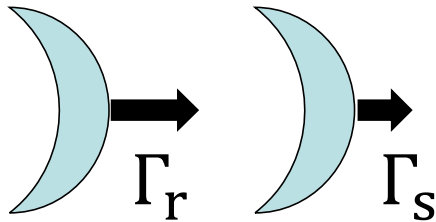
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## Prompt Emission

- ***Eiso =  $10^{52} - 10^{55}$  erg may require high efficiency of energy dissipation.***
- ***Several gamma-ray polarization reports imply synchrotron emission. (Some can be photospheric)***
- ***Jets are driven by Blandford-Znajek process?***
- ***Then, magnetically dominated jets?***
- ***The classical internal shock model works?***

# Objection to Internal Shock Model

## Classical picture



collide of two shells → merge

Shocked Region			
1	2	3	4
→ $\Gamma_{12}$ $n'_r$	$U'_2 = (\Gamma_{12} - 1)n'_2 m_p c^2$ $n'_2 = (4\Gamma_{12} + 3)n'_r$	$U'_3 = U'_2$ $n'_3 = (4\Gamma_{34} + 3)n'_s$	← $\Gamma_{34}$ $n'_s$

## Energy dissipation efficiency

$$\Gamma_m \approx \sqrt{\frac{M_r \Gamma_r + M_s \Gamma_s}{M_r / \Gamma_r + M_s / \Gamma_s}}$$

$$f \equiv \frac{\Gamma_m E_{\text{int}}}{(M_r \Gamma_r + M_s \Gamma_s) c^2} = 1 - \frac{2\sqrt{\Gamma_r / \Gamma_s}}{1 + \Gamma_r / \Gamma_s} \approx 1 - 2\sqrt{\Gamma_s / \Gamma_r}$$

↑ Equal mass      ↑  $\Gamma_r \gg \Gamma_s$

**If**  $\Gamma_r / \Gamma_s = 2 \Rightarrow f = 0.057$   
 $\Gamma_r / \Gamma_s = 10 \Rightarrow f = 0.43$

**Large dispersion in  $\Gamma$  is required for high  $f$ .**

## Objection to Internal Shock Model

**Typical Synchrotron Photon Energy**  $\varepsilon_{pk} = \frac{\hbar e B}{m_e c} \gamma_m^2 \Gamma_m$

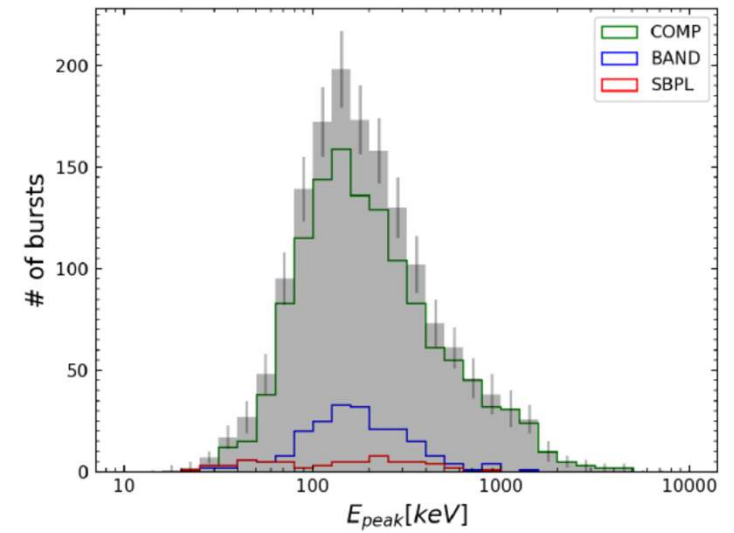
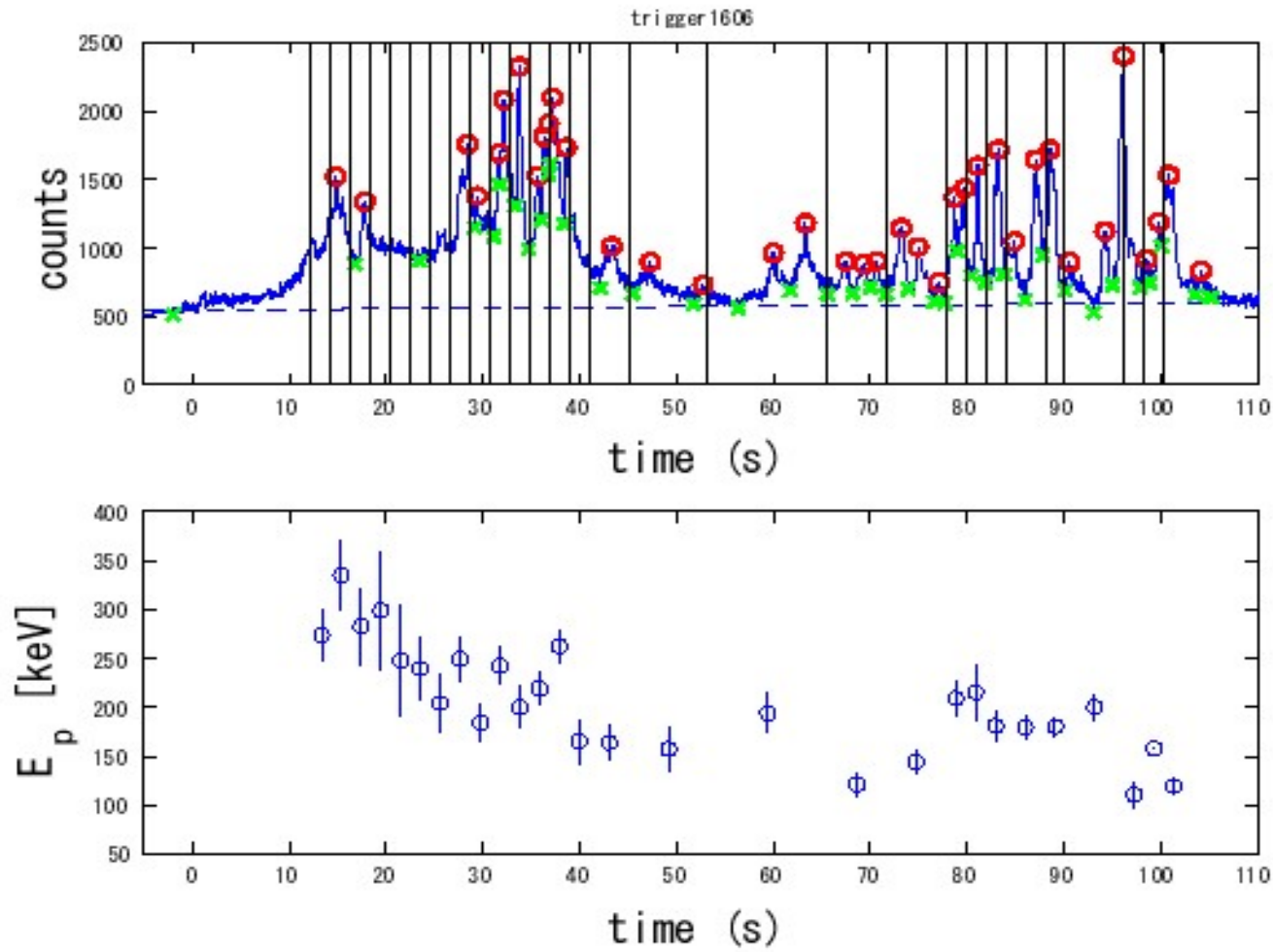
$$\frac{\Gamma_r}{\Gamma_s} = 10 \rightarrow \Gamma_{sh} = 2.57$$

$$\frac{\Gamma_r}{\Gamma_s} = 2 \rightarrow \Gamma_{sh} = 1.08$$

$$\gamma_m \propto \epsilon_e (\Gamma_{sh} - 1) \frac{m_p}{m_e} \rightarrow \varepsilon_{pk} \propto (\Gamma_{sh} - 1)^2$$

**For large dispersion in  $\Gamma$ , large dispersion in  $\varepsilon_{pk}$**

## Small Dispersion in typical photon energy



**Fermi GBM GRB Spectral Catalog 21**

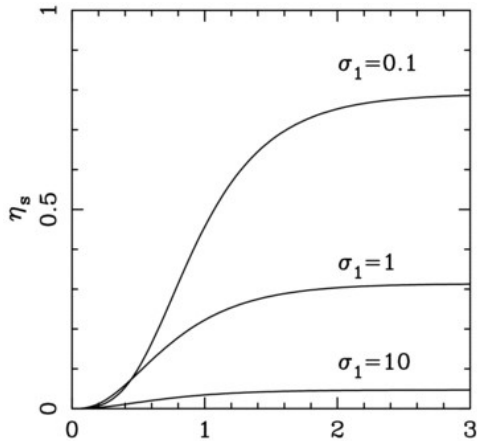
# Magnetically dominated case

$$\sigma \equiv \frac{B^2}{4\pi\rho c^2}$$

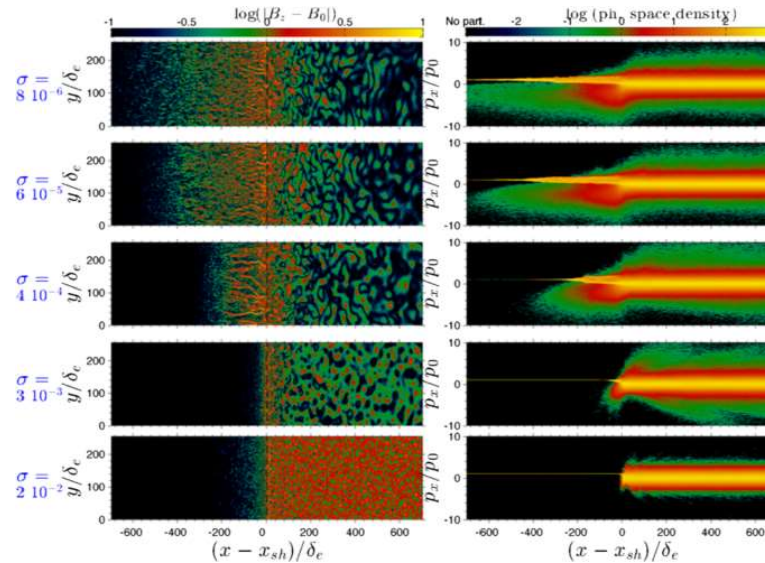
- **High-sigma case**,  $\Gamma_{sh} \approx \sqrt{\sigma}$ .
- **Energy dissipation due to shocks becomes small.**

**Thermal energy fraction in the downstream  $\frac{3}{8\sigma}$**

## Dissipation Efficiency



Komissarov 12



Plotnikov + 18

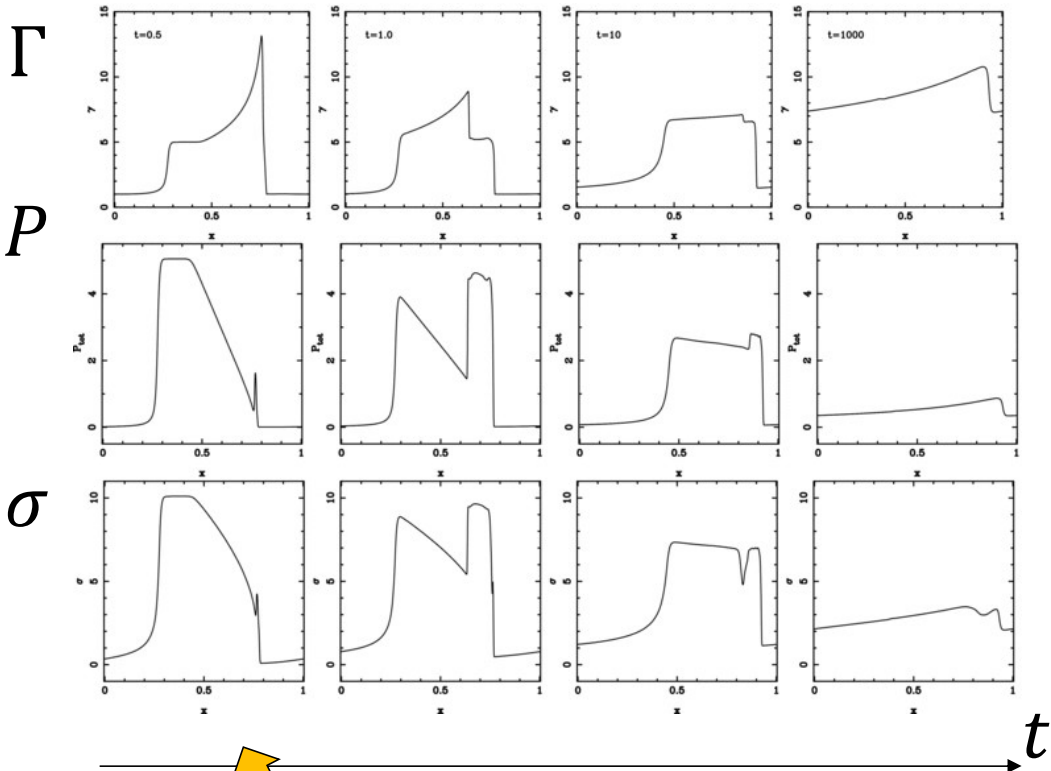
**In addition, shock acceleration is hard for high sigma.**

**Need Reconnection?**

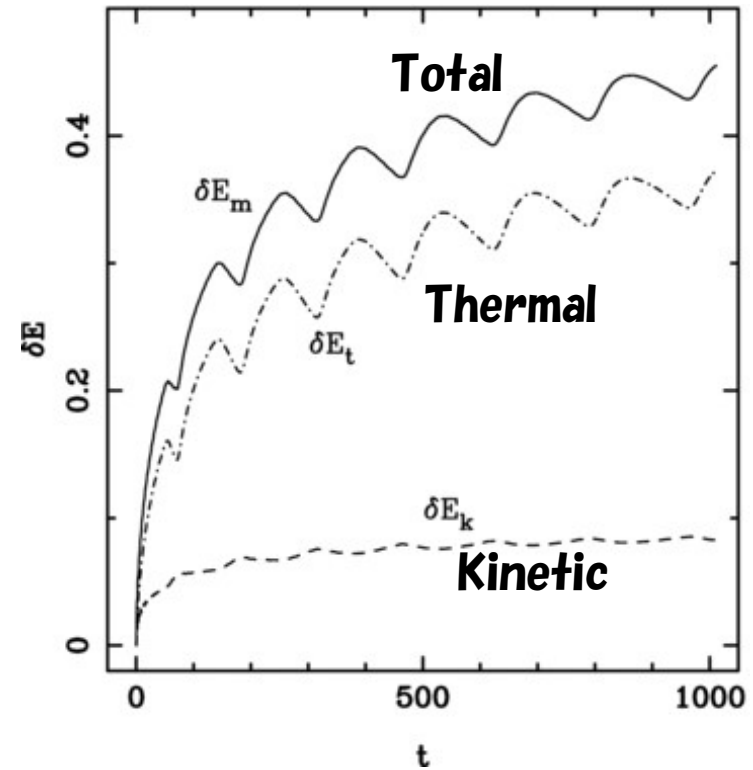
# Shock dissipation of high-sigma flows

Komissarov 12

Periodic Boundary



Impulsive acceleration see Granot+ 11

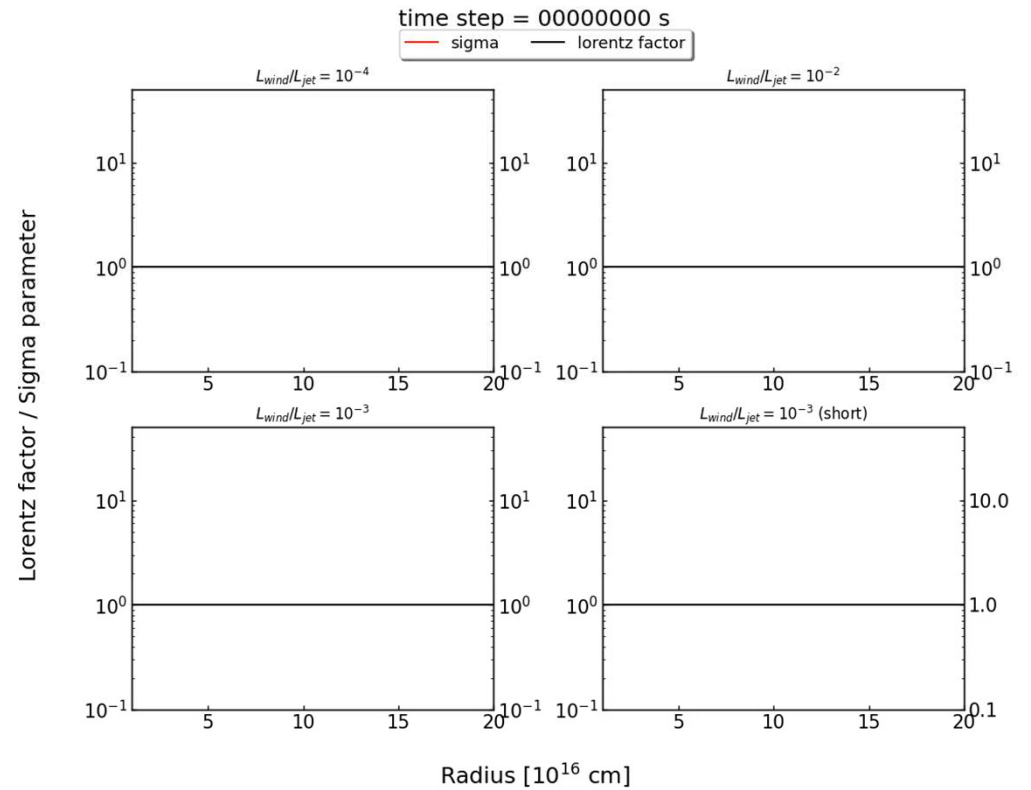
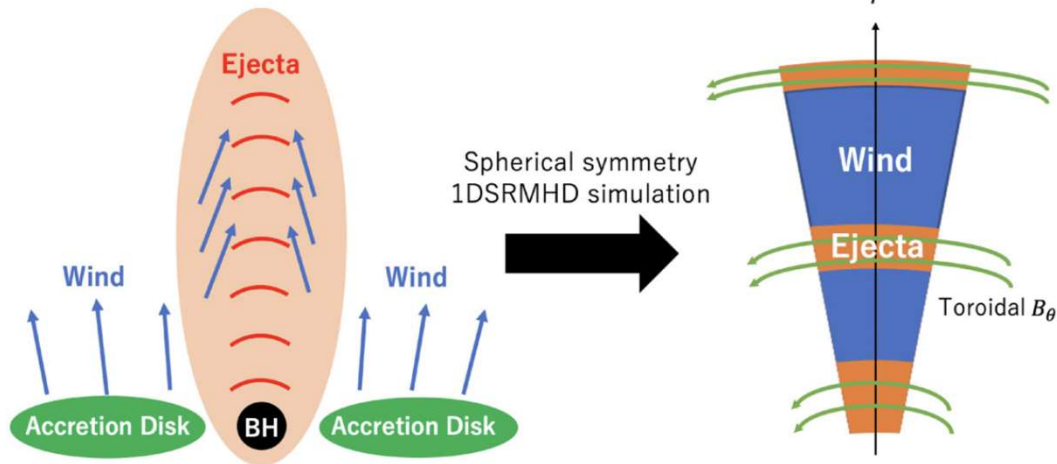


Dissipated Magnetic Energy

# 1D-spherical simulations

Kusafuka, KA+ 23

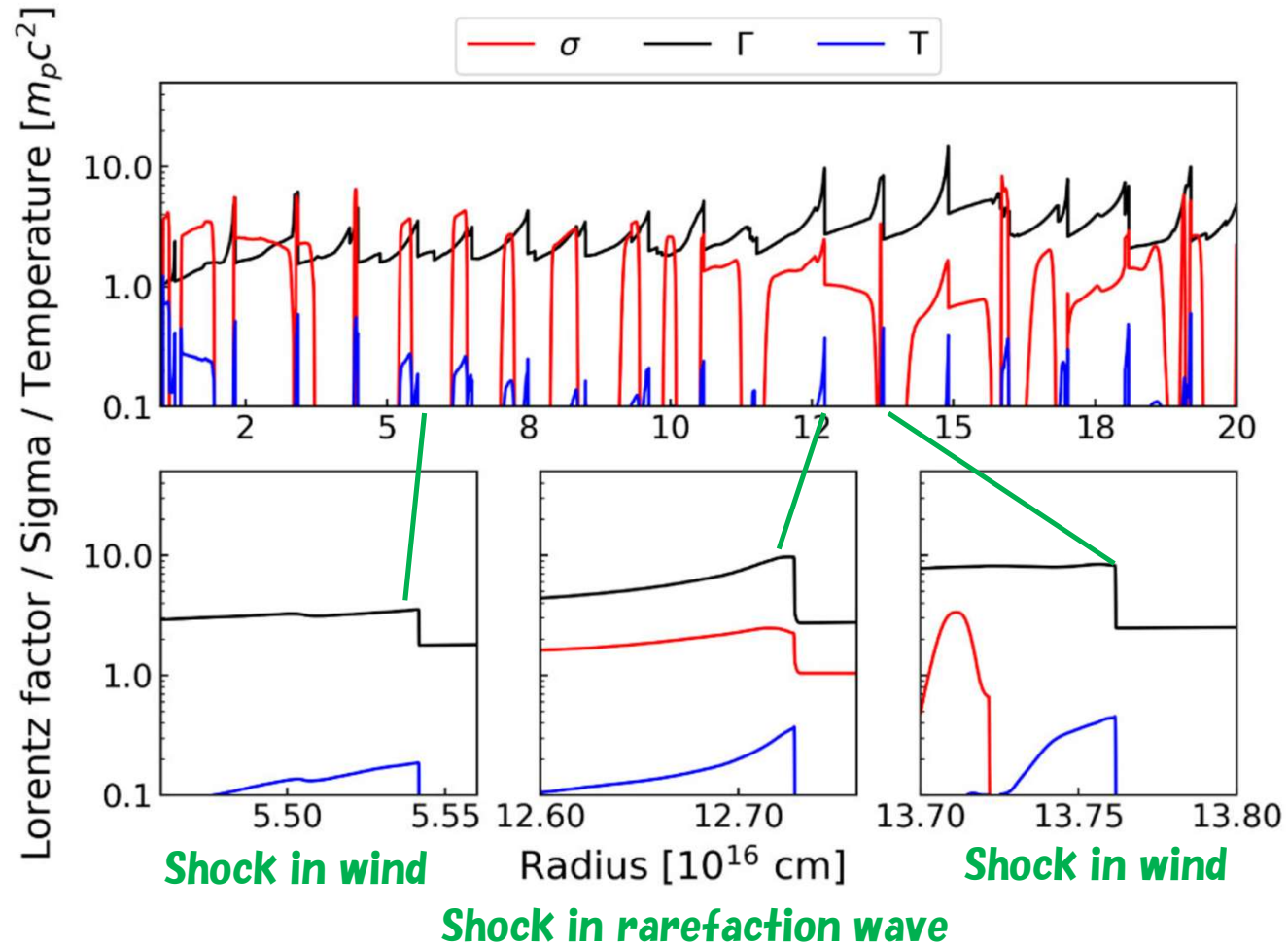
## Intermittent ejection





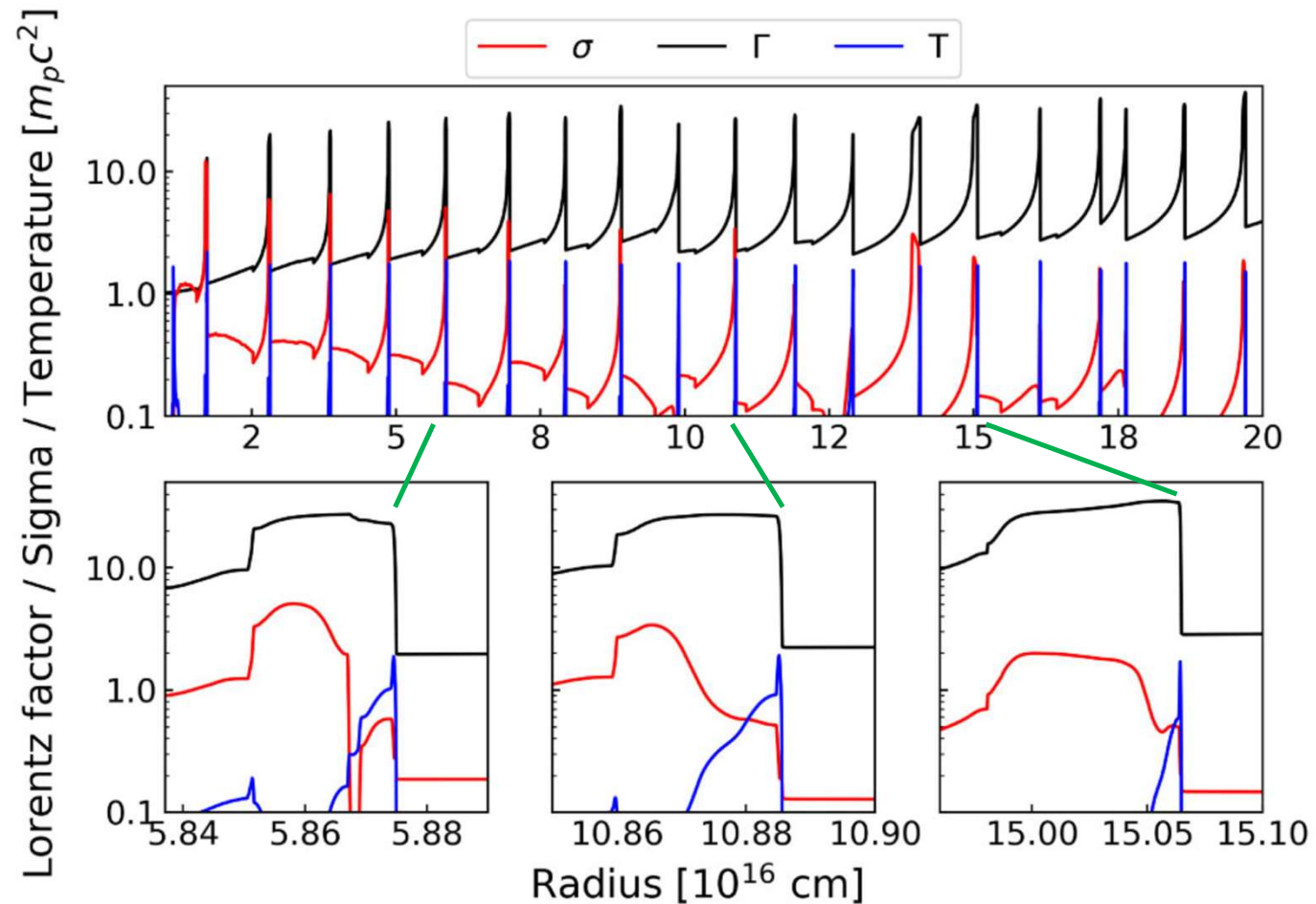
# Shocks in wind

$$L_w = 10^{-2} L_j$$



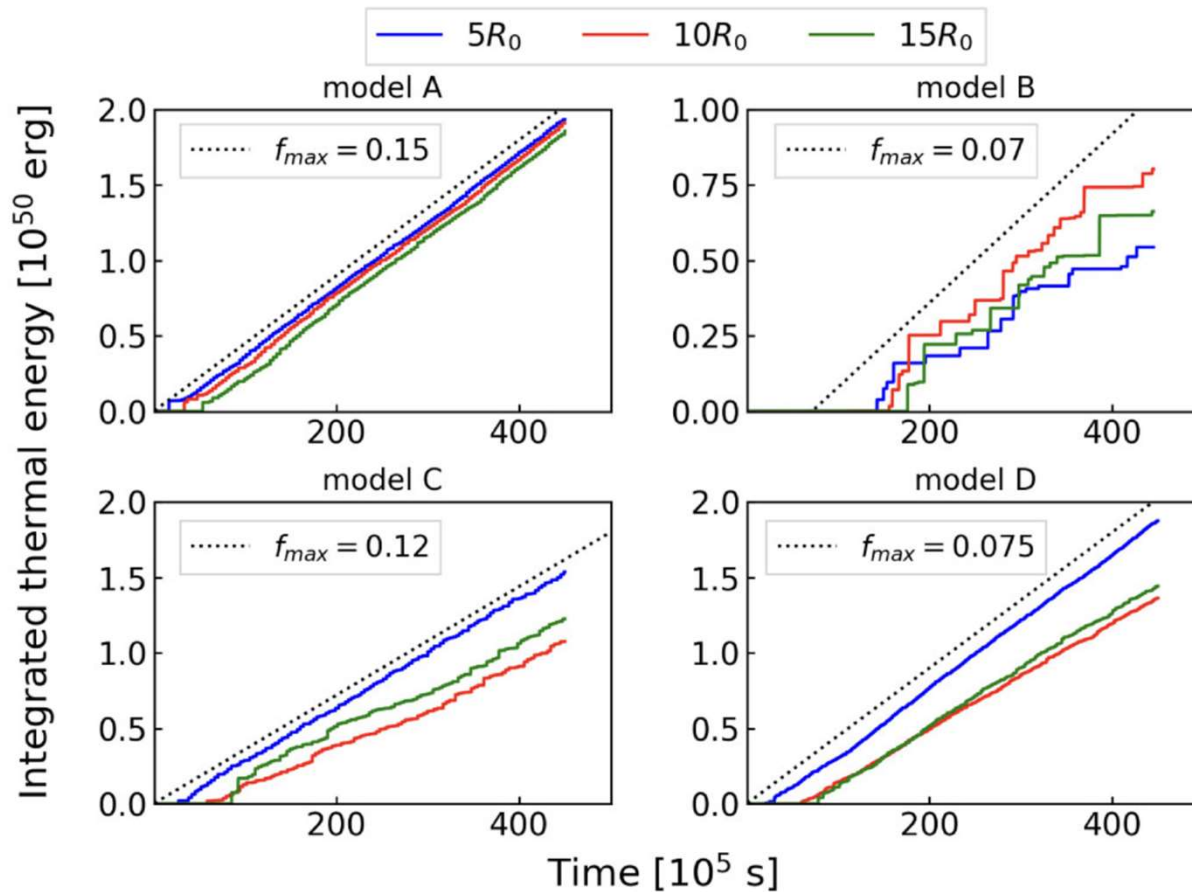
## Shocks in rarefaction wave

$$L_w = 10^{-4} L_j$$



**Jet compresses wind to be crushed.**

# Energy Dissipation



$\sim 10\%$

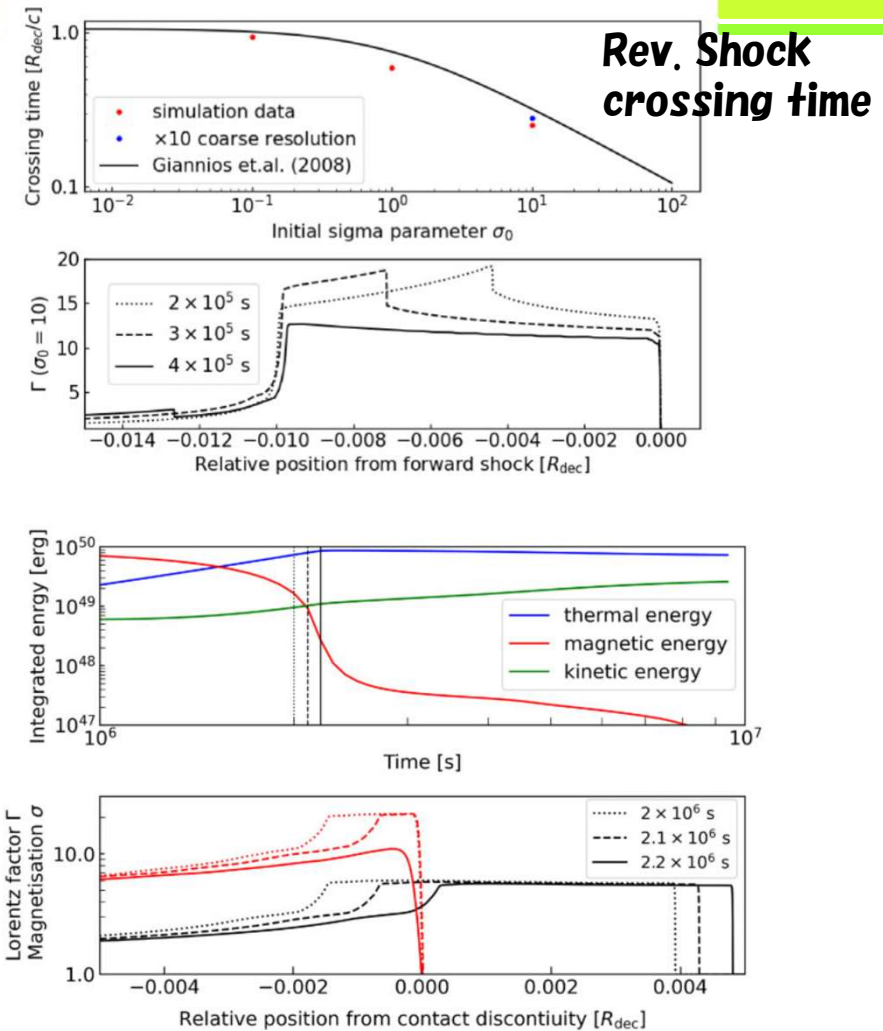
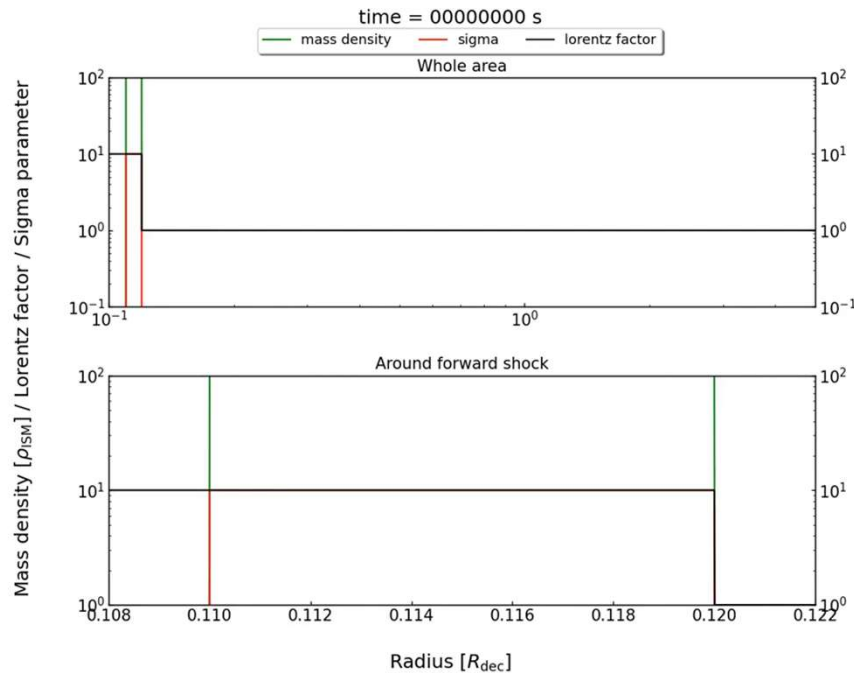
**No need for  
large dispersion in  $\Gamma$**

**Moderate dispersion in  $\Gamma_{sh}$   
→ small dispersion in  $\varepsilon_{pk}$**

**$\Gamma > 100$  in GRBs,  
which may enhance efficiency**

# External Shock (Sorry despite session for prompt emission ...)

- External shock is also a way to dissipate the magnetic energy.

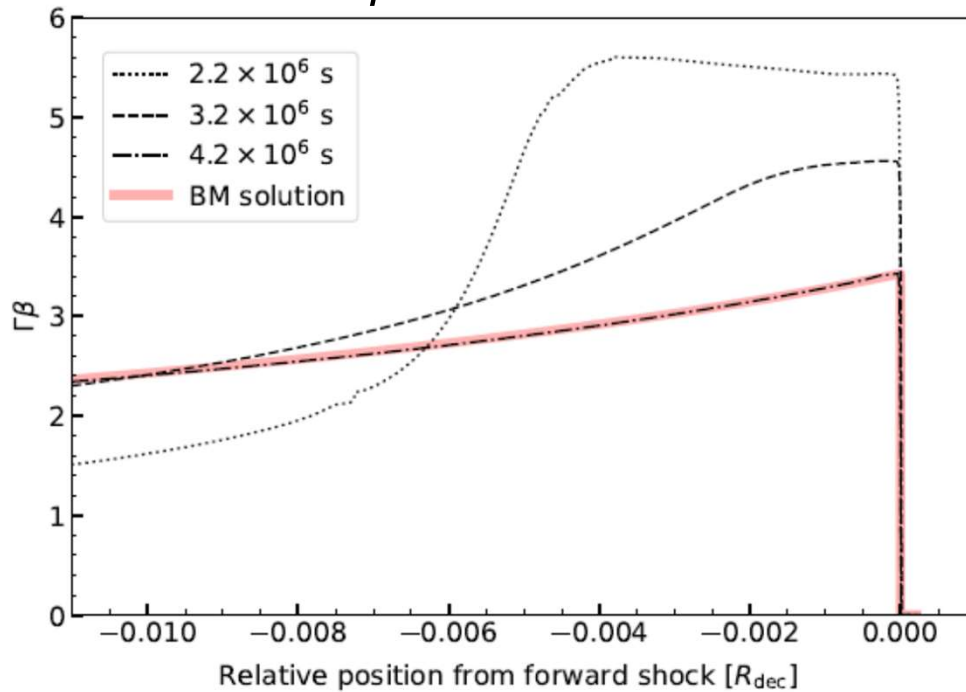


# Transition Phase=Energy Injection Phase

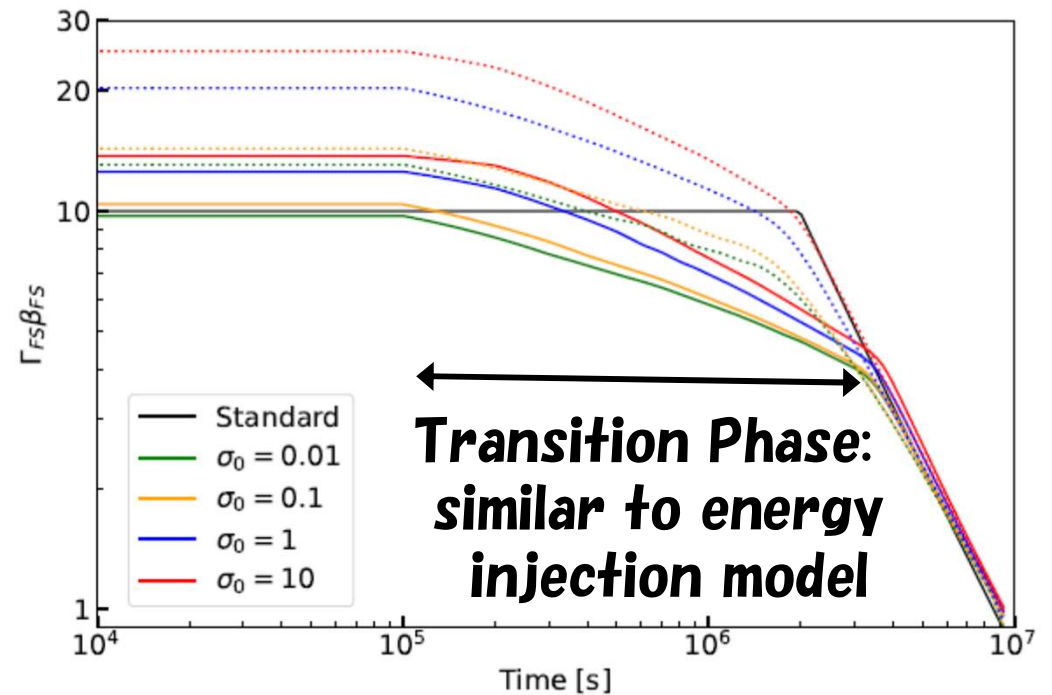
Kusafuka & KA 24

See also Kobayashi & Sari 2000; Mimica et al. 2009

## Evolution of $\Gamma\beta$ -Profile

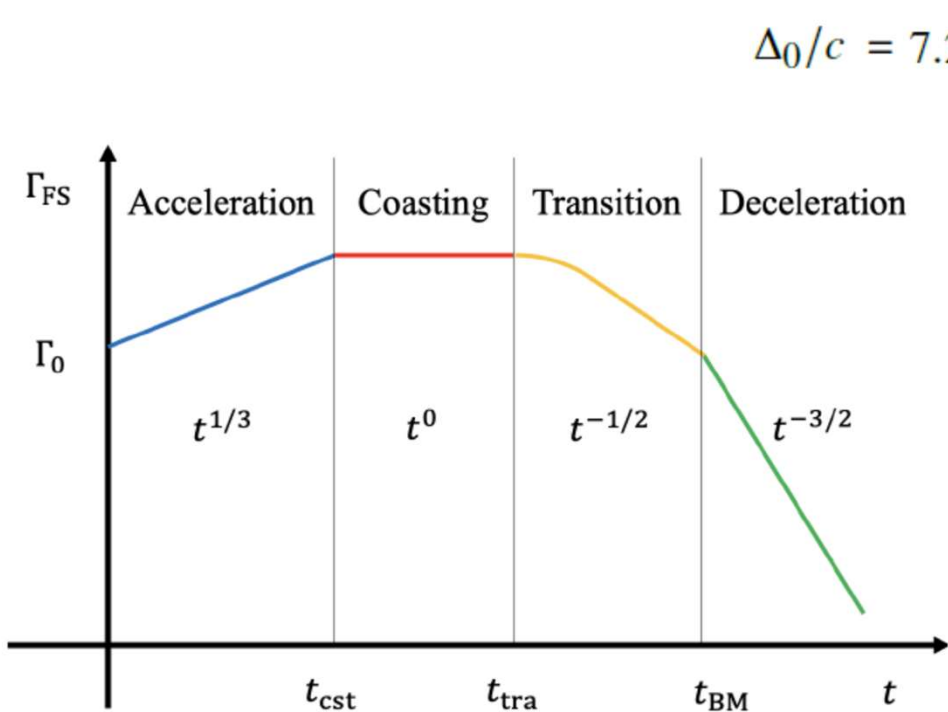


## Forward Shock $\Gamma\beta$

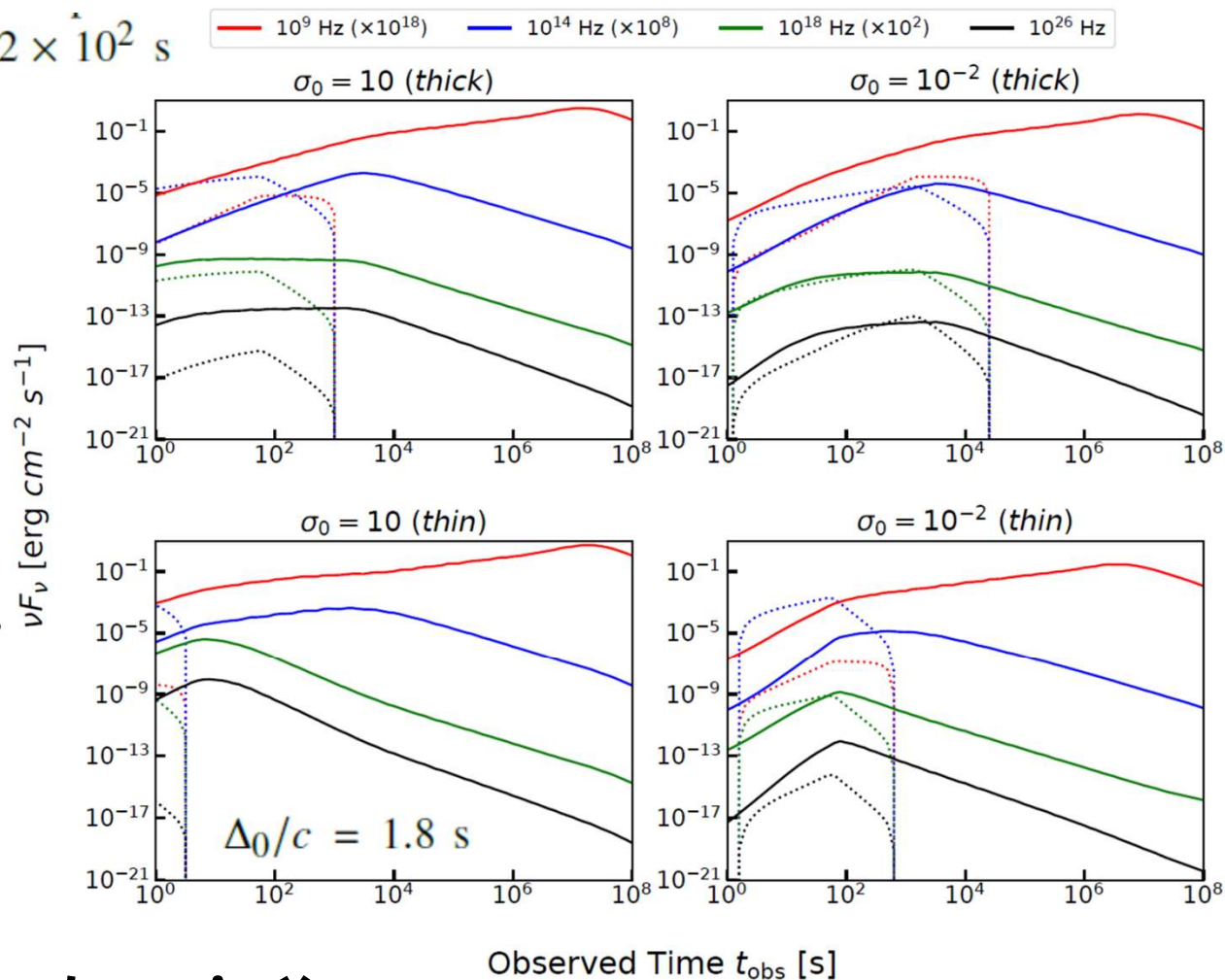


**Finite width of ejecta: delay of transition to BM solution.  
The transition phase lasts longer than the reverse shock crossing time.**

# Afterglow lightcurve



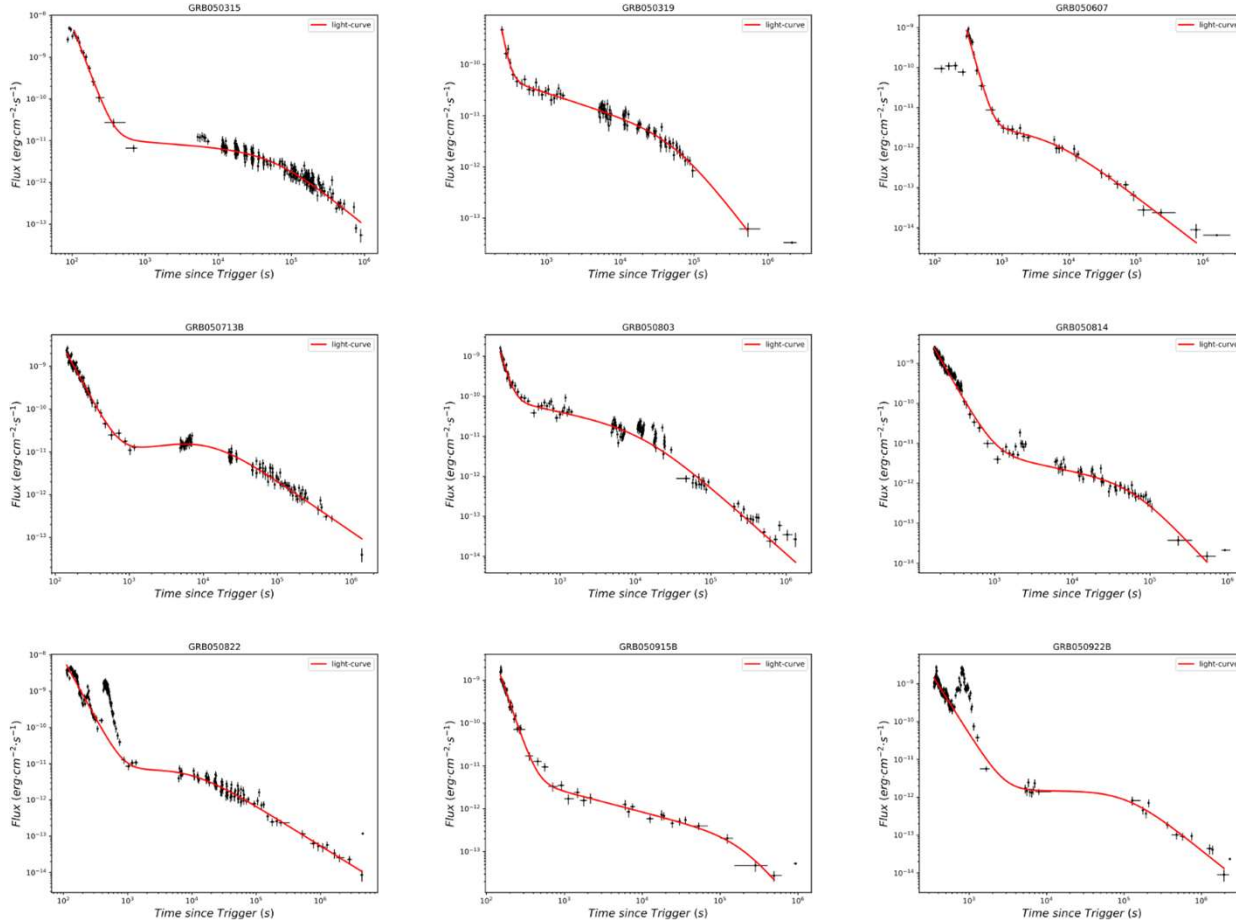
$E_0 = 10^{52} \text{ erg}, \Gamma_0 = 100, \text{ and } n_0 = 1 \text{ cm}^{-3}.$   
 $\epsilon_e = 0.1 \quad \epsilon_B = 0.01 \quad p = 2.2$



**Transition phase leads to shallow decay in X-ray**

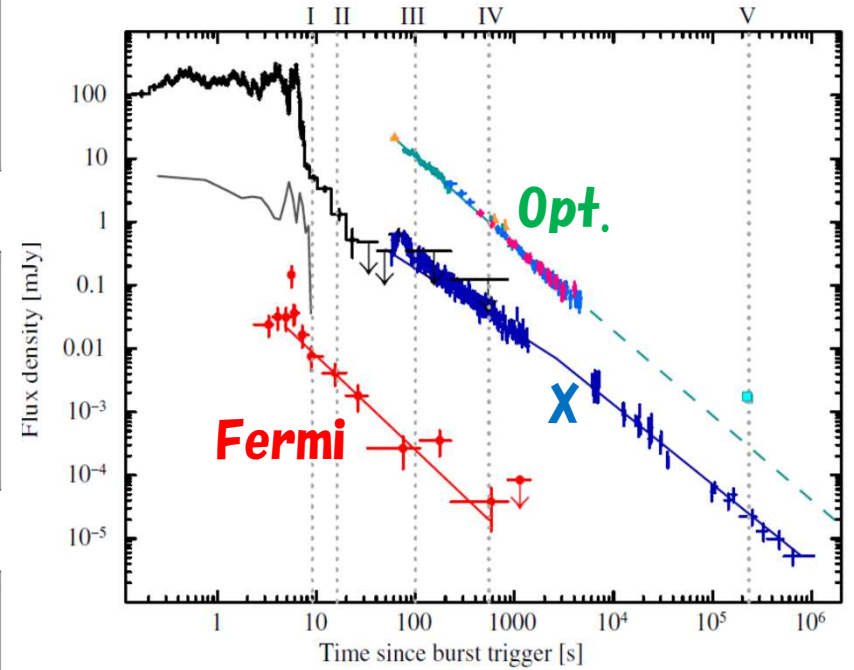
# Shallow Decay Phase

## X-ray LCs



X.-K. Ding+ 22

## GRB 110731A (Ackerman+ 13)



Fermi-LAT GRBs tend to show no shallow decay.

See Yamazaki+ 20

# 1D time-dependent calculation

Fukushima, KA+ 2017

**Mass:**

$$\frac{dM}{dt} = \frac{1}{\Gamma} \frac{dM}{dt'} = 4\pi R^2 c \beta_{\text{sh}} n_{\text{ISM}} m_p,$$

**Energy (Lorentz factor):**

$$\frac{dE'_{\text{sh}}}{dt'} = \Gamma c^2 \frac{dM}{dt'} - \frac{dE'_{\text{rad}}}{dt'} - \frac{dE'_{\text{ad}}}{dt'},$$

$$E_{\text{sh}} = \Gamma E'_{\text{sh}} = E_0 + Mc^2 - E_{\text{rad}}.$$

**Density and Volume:**

$$n' = \frac{\hat{\gamma}\Gamma + 1}{\hat{\gamma} - 1} n_{\text{ISM}},$$

$$V' = M / (m_p n').$$

**Magnetic Field:**

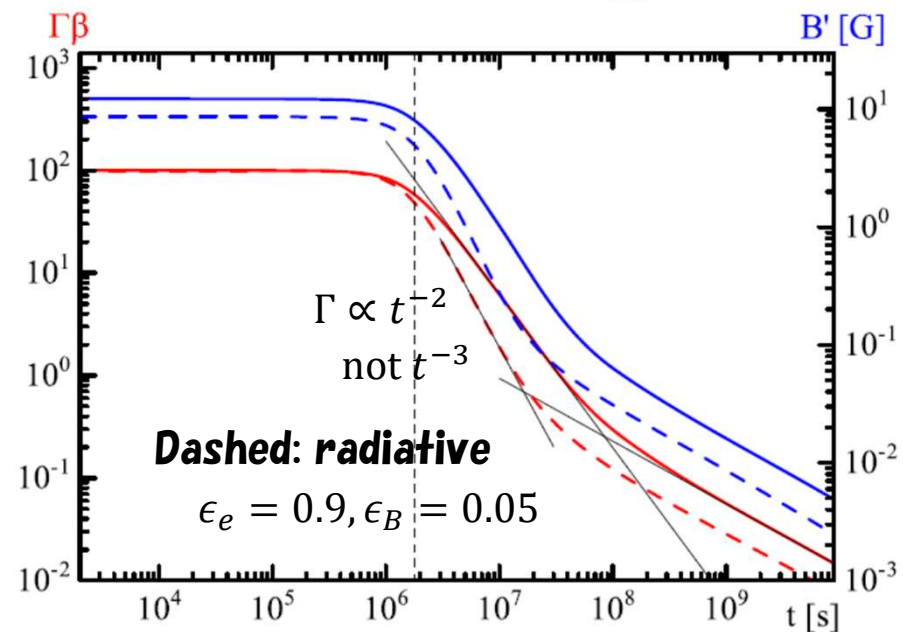
$$\frac{dE'_B}{dt'} = \epsilon_B (\Gamma - 1) c^2 \frac{dM}{dt'}.$$

$$B' = \sqrt{\frac{8\pi E'_B}{V'}}.$$

**Non-thermal electrons:**

$$\frac{dN'_e}{dt'} = \frac{\eta}{m_p} \frac{dM}{dt'} = \int_{\epsilon'_{\text{min}}}^{\infty} d\epsilon' \dot{N}'_e(\epsilon'_e),$$

$$\frac{dE'_e}{dt'} = \epsilon_e (\Gamma - 1) c^2 \frac{dM}{dt'} = \int_{\epsilon'_{\text{min}}}^{\infty} d\epsilon' \epsilon'_e \dot{N}'_e(\epsilon'_e).$$

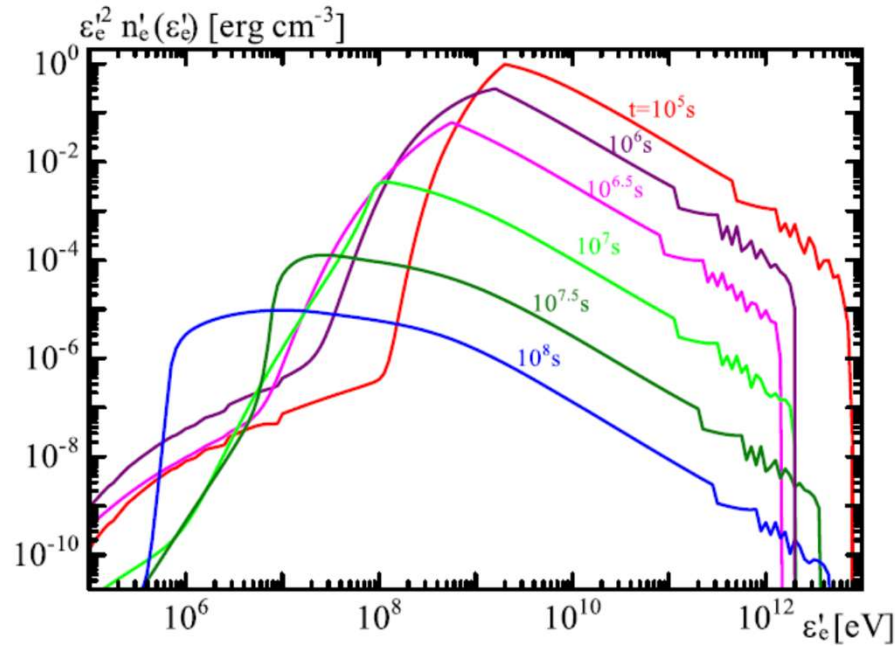


$$E_0 = 10^{52} \text{ erg}, \Gamma_0 = 100, n_{\text{ISM}} = 1 \text{ cm}^{-3}, p = 2.2, \\ \epsilon_e = \epsilon_B = 0.1$$



# 1D time-dependent calculation

## Electron energy distribution



以下では基本的に

$$E_0 = 10^{53} \text{ erg}, \epsilon_e = 0.1, \epsilon_B = 10^{-3}$$

$$\Gamma_0 = 300, n_0 = 1 \text{ cm}^{-3}, \text{ and } p = 2.2$$

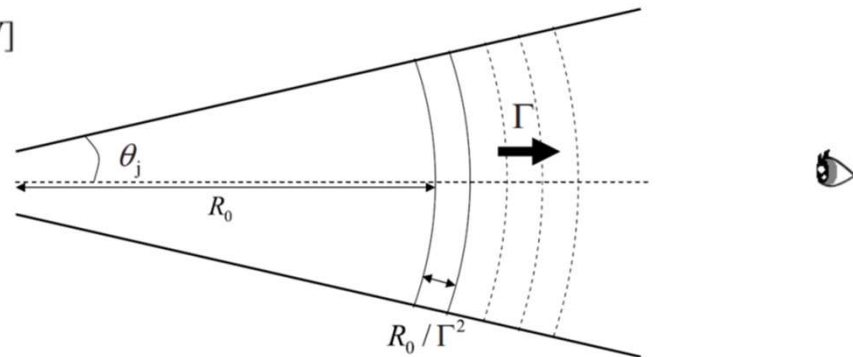
Taking into account the theta-dependence on observables.

$$\epsilon_{\text{obs}} = \frac{\epsilon'}{\Gamma(1 - \beta \cos \theta)(1 + z)},$$

$$t_{\text{obs}} = (1 + z) \left[ \left( t - \frac{R - R_0}{c} \cos \theta \right) + \frac{R_0}{c} (1 - \cos \theta) \right]$$

$$\Phi(\epsilon_{\text{obs}}, t_{\text{obs}}) = \frac{dN_\gamma}{dS_{\text{obs}} dt_{\text{obs}} d\epsilon_{\text{obs}}}$$

$$= \int d\theta \left( \frac{R(t'_\theta)}{D} \right)^2 \frac{\sin \theta |\cos \theta'| c n'_\gamma(\epsilon'_\theta, t'_\theta)}{2\Gamma^2(1 - \beta \cos \theta)(1 - \beta_{\text{sh}} \cos \theta)},$$



# Energy Injection Model

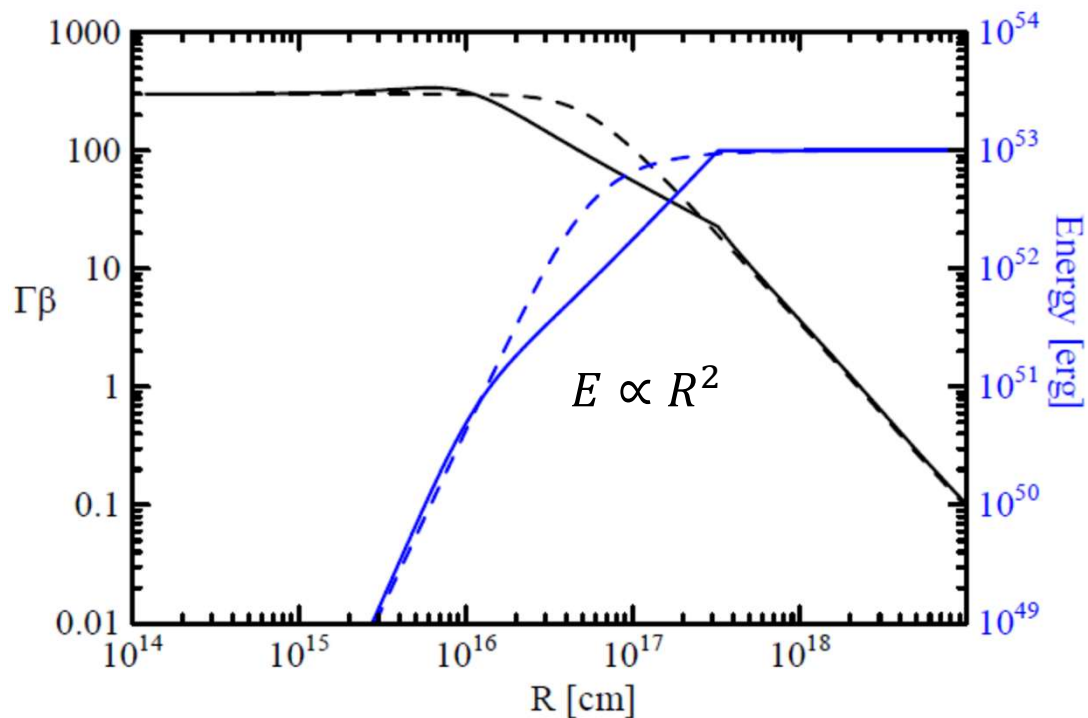
Late Jet, catching up

or transition phase  
(phenomenologically equivalent)

Shock Front

$$E = E_i + (E_0 - E_i) \left( \frac{R}{R_e} \right)^2, \quad \text{for } R < R_e.$$

$$R_e = 3.2 \times 10^{17} \text{ cm} \quad E_i = 10^{51} \text{ erg}$$



$$E \propto \bar{R}^2 \text{ leads to } \Gamma \propto R^{-1/2}$$

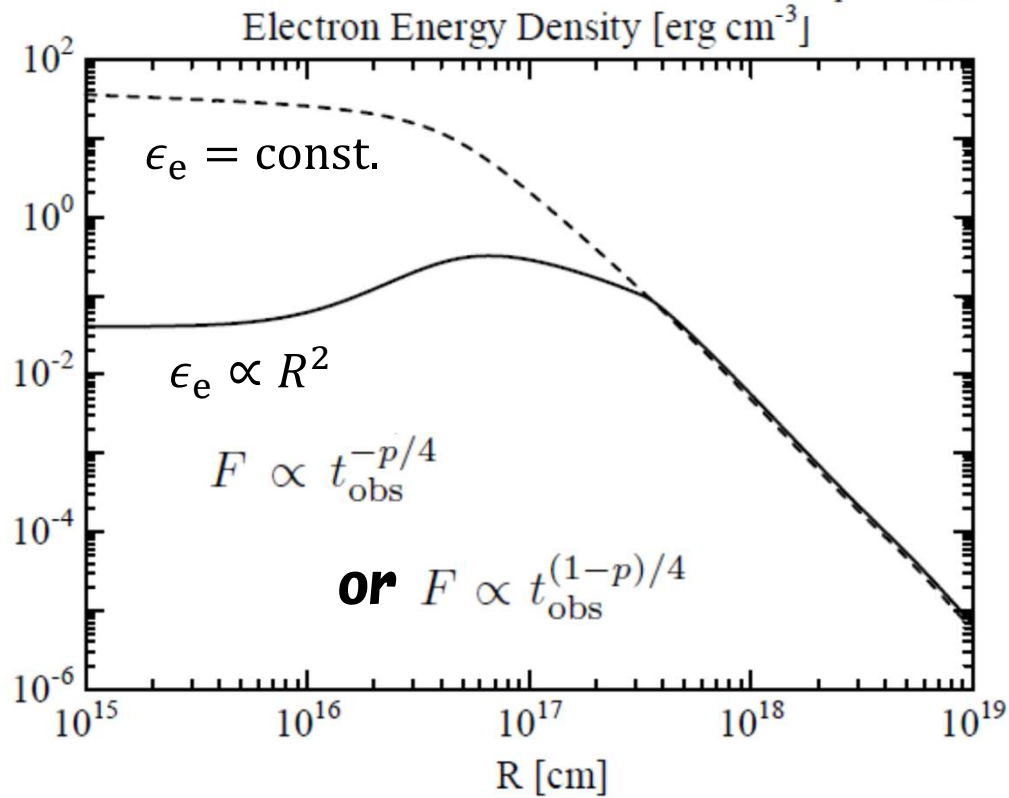
**but**  $\Gamma \propto R^{-0.8}$  for  $R < R_e$ .

# Phenomenological Model

*Evolving  $\epsilon_e$*

$$\epsilon_e = \epsilon_i + (\epsilon_f - \epsilon_i) \left( \frac{R}{R_e} \right)^2, \quad \text{for } R < R_e.$$

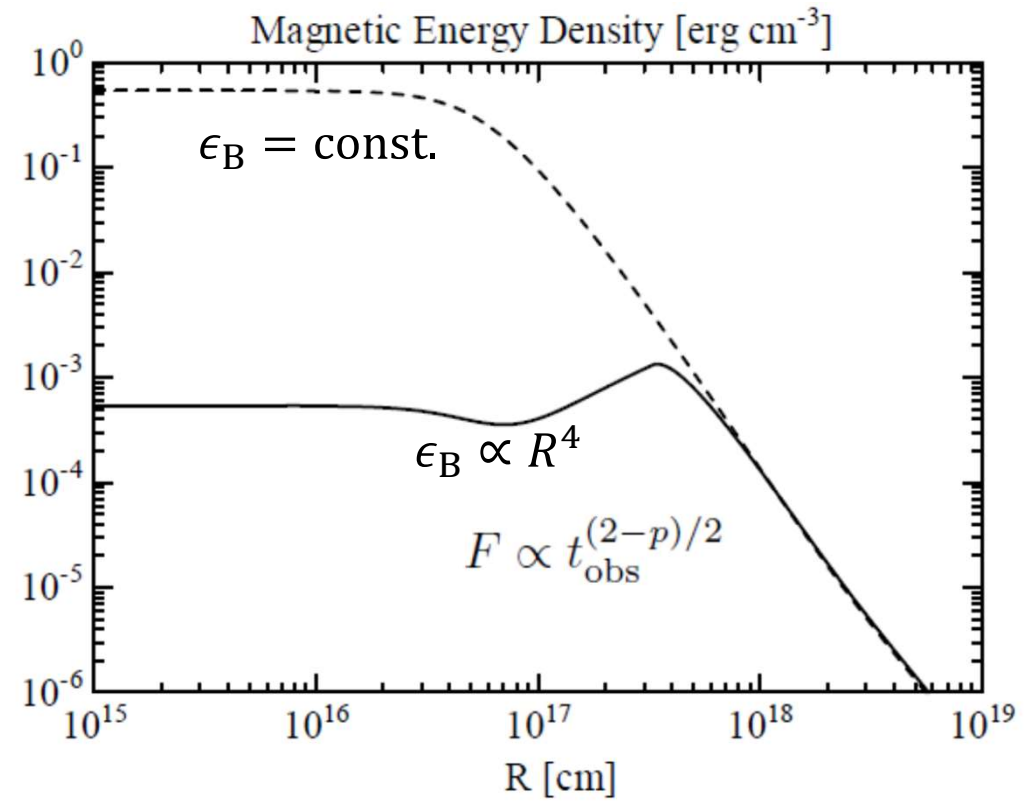
$\epsilon_i = 10^{-4}$



*Evolving  $\epsilon_B$*

$$\epsilon_B = \epsilon_i + (\epsilon_f - \epsilon_i) \left( \frac{R}{R_e} \right)^4, \quad \text{for } R < R_e.$$

$\epsilon_i = 10^{-6}$



# Low- $\Gamma$ + Wind

$$E_0 = \Gamma_0^2 M(R_{\text{dec}}) c^2$$

$$t_{\text{dec}} \simeq (1+z) \frac{R_{\text{dec}}}{5c\Gamma_0^2}$$

$$\simeq 150(1+z) A_*^{-1} \left( \frac{E_0}{10^{53} \text{ erg}} \right) \left( \frac{\Gamma_0}{30} \right)^{-4} \text{ s}$$

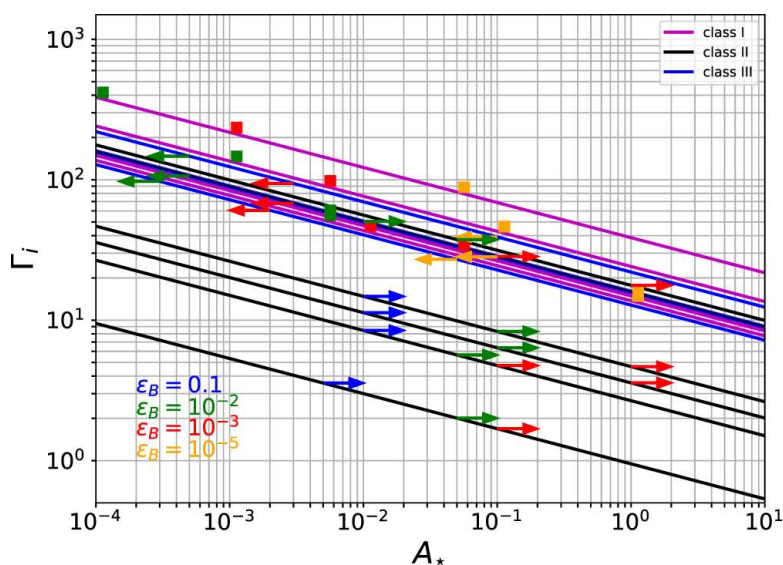
coasting phase,  $F \propto t_{\text{obs}}^{(2-p)/2}$

$A_* = 0.3$ ,  $E_0 = 2 \times 10^{53}$  erg, and  $\Gamma_0 = 30$

$t_{\text{dec}} \simeq 1500$  s for  $z = 0.5$ .

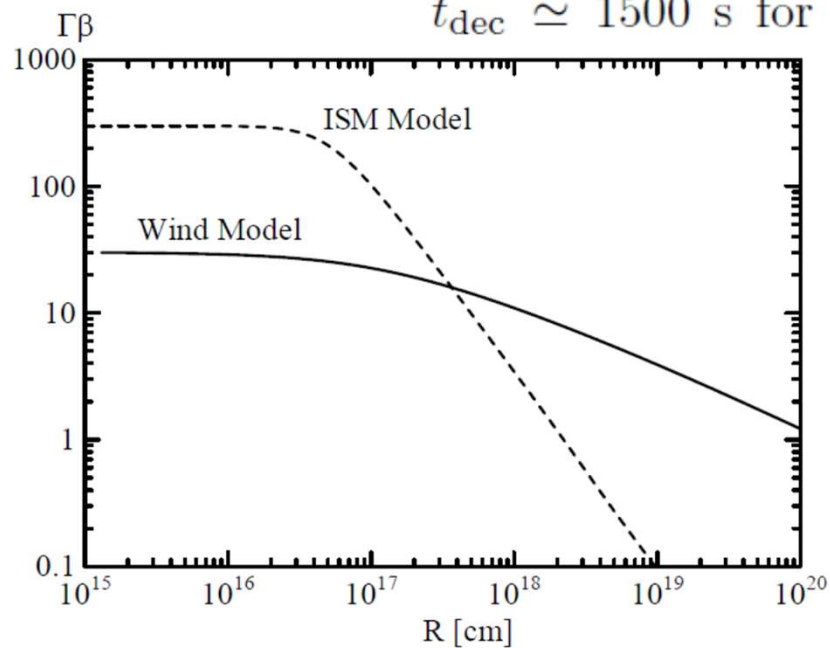
Wind Model

See Dereli-Bégué +22



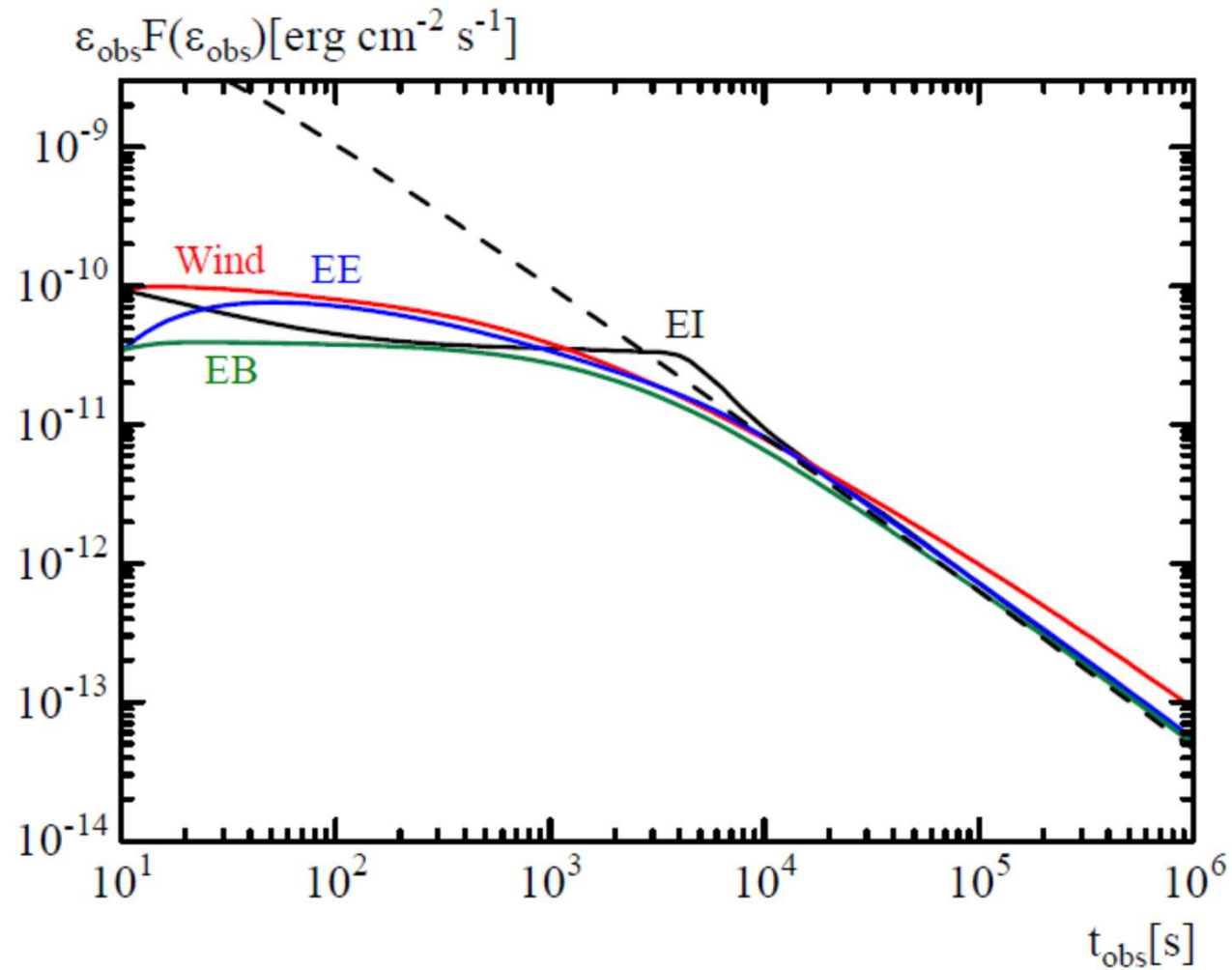
$$\rho = 5.0 \times 10^{11} A_* \left( \frac{R}{1 \text{ cm}} \right)^{-2} \text{ g cm}^{-3}$$

$A_* = 1$  corresponds to  $10^{-5} M_{\odot} \text{ yr}^{-1}$

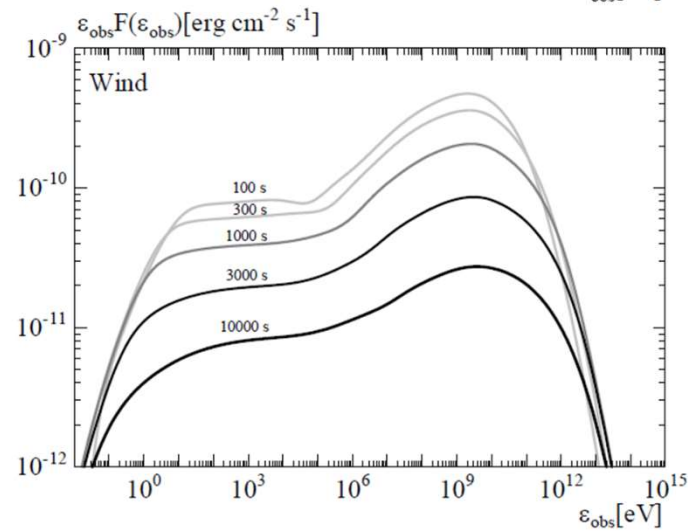
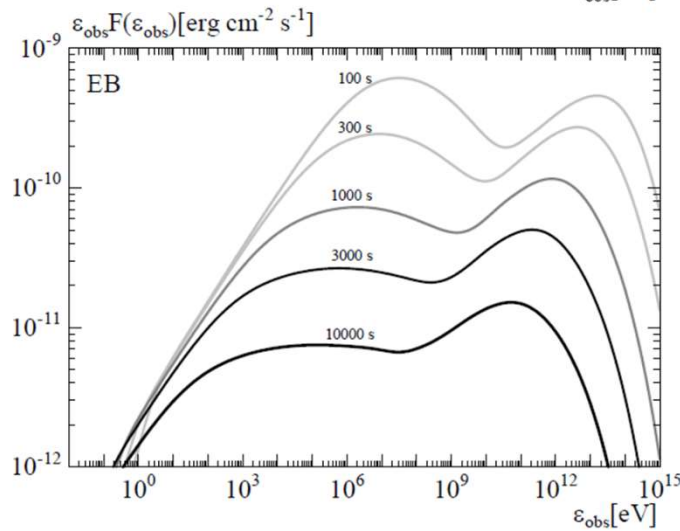
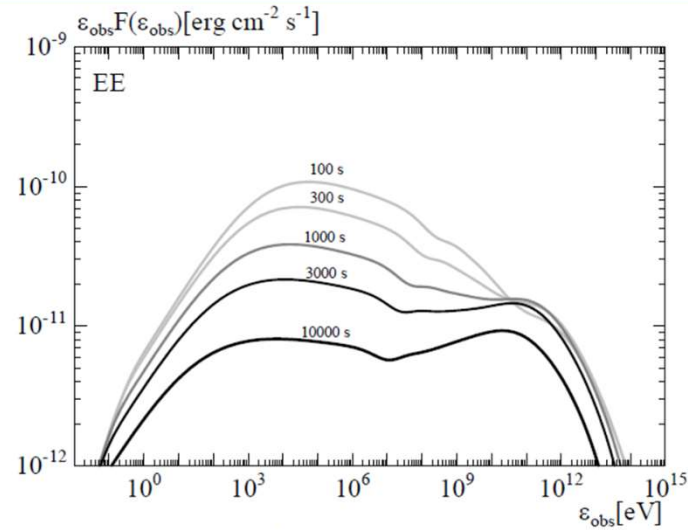
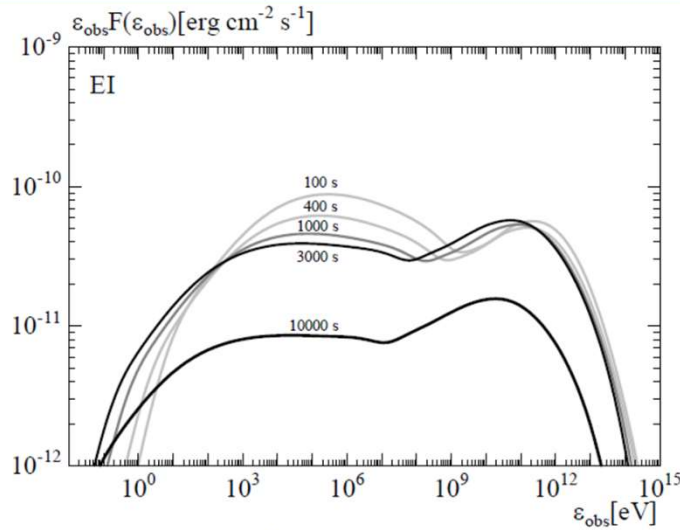


# Lightcurves

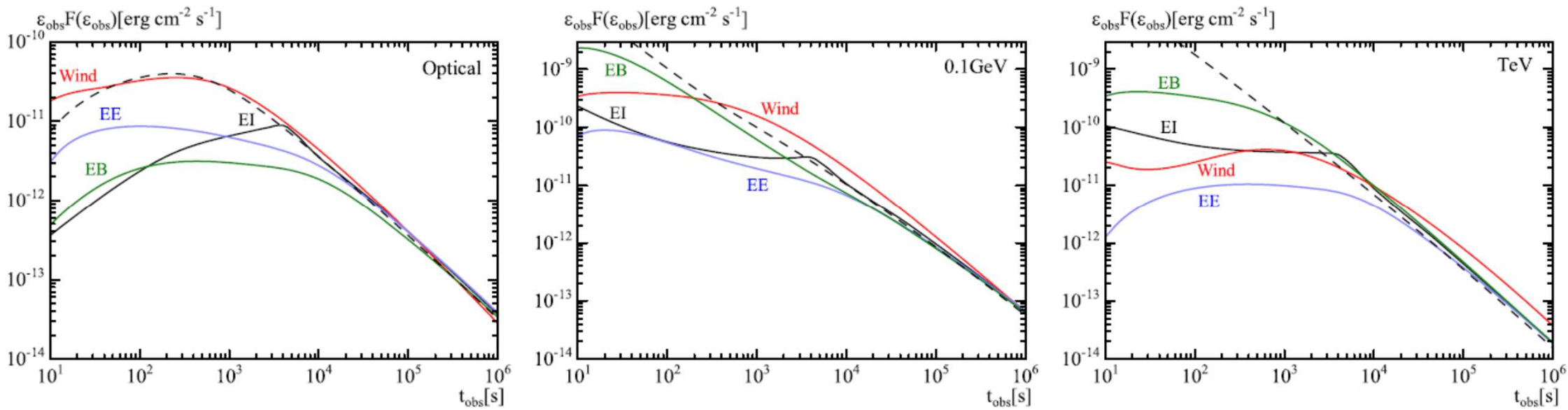
X-ray (1keV)



# Spectrum Evolution



# Multi-wavelength Lightcurves



***Optical can be polluted by reverse shock component.  
Variety in TeV lightcurves.***

## Summary

- **Internal Shock**. High sigma outflow can be dissipated via interaction with winds or rarefaction waves.
- **External Shock**. Transition phase before catch-up of rarefaction wave = Energy injection phase. Relevant for Shallow-decay phase?
- **GRB 240529A**. Shallow decay with gamma-rays.
- **BOAT GRB**. Model without early jet break